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Structure of Hydridochlorobis(diphenylethylphosphine)platinum¹

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The crystal and molecular structure of hydridochlorobis(diphenylethylphosphine)platinum, $\text{PtHCl}[\text{P}(\text{C}_6\text{H}_5)_2\text{C}_2\text{H}_5]_2$, has been determined from three-dimensional X-ray data. The material crystallizes in space group $\text{C}_{2h}^2\text{-P2}_1/\text{c}$ of the monoclinic system, with four molecules in a cell of dimensions $a = 11.80$, $b = 16.93$, $c = 14.31$ Å., $\beta = 108.4^\circ$. The individual monomeric molecules are well separated, the closest Pt-Pt distance being greater than 9 Å. The phosphorus atoms are *trans* to one another and with the chlorine lie at three of the four corners of an approximate square, with the platinum in the center. In agreement with results reported previously for a similar bromo compound, the Pt-P distances are shorter and the Pt-halogen distance longer than are calculated from the usual radii-sum rules. The Pt-halogen lengthening is consistent with the high lability of the halogen *trans* to a hydrogen. The platinum, phosphorus, and chlorine atoms deviate slightly, but significantly, from coplanarity.

Introduction

The nature of the metal-hydrogen bond in transition metal hydrides has received renewed attention and re-interpretation in the past few years, primarily as the result of the discovery of new classes of stable transition metal hydrides³ and because diffraction studies⁴⁻⁹ have shown that in such compounds the hydrogen exerts an important influence on the stereochemistry. This influence was not apparent in earlier chemical^{10,11} and spectroscopic studies.¹²⁻¹⁴ Of paramount importance in this re-interpretation of the metal-hydrogen bond was the determination by Owston, Partridge, and Rowe⁴ of the structure of hydrido-bromobis(triethylphosphine)platinum. They showed that the *trans* P and Br atoms are at three of the four corners of a square, with the Pt at the center. They

found the Pt-P distances to be shorter and the Pt-Br distance to be longer than those predicted from radii-sum rules. Although they obtained no direct evidence for the hydrogen atom location, it seems clear that hydrogen is occupying the fourth corner of the square, for otherwise there would be no reason for the other atoms to be situated at three of the four corners of a square. Their determination of the structure was hampered considerably by the X-ray sensitivity of the crystals, and accordingly it was based on minimal zonal data. Because of this, the limits of error are very high, being ± 0.07 Å. for Pt-P and ± 0.04 Å. for Pt-Br. The light atoms were not located with any accuracy.

Because of the central role played by this structure determination in the revision of ideas concerning the stereochemical influence of hydrogen in transition metal hydrides, a determination to much higher accuracy on a similar compound seems worthwhile. In addition it is important to establish the bond distances with higher accuracy for use in an eventual description of the bonding and there is the possibility, though faint, that the hydride hydrogen might be located directly. In this report on the structure of hydridochlorobis(diphenylethylphosphine)platinum we indeed do define the geometrical arrangement with much higher accuracy, but we fail to provide direct evidence concerning the hydrogen position.

Collection and Reduction of the X-Ray Data

The white crystals of $\text{PtHCl}[\text{P}(\text{C}_6\text{H}_5)_2\text{C}_2\text{H}_5]_2$ were prepared by H. B. Gray using the method described

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(3) See, for example, J. Chatt, *Proc. Chem. Soc.*, 318 (1962).

(4) P. G. Owston, J. M. Partridge, and J. M. Rowe, *Acta Cryst.*, **13**, 246 (1960).

(5) P. L. Orioli and L. Vaska, *Proc. Chem. Soc.*, 333 (1962).

(6) S. J. La Placa and J. A. Ibers, *J. Am. Chem. Soc.*, **85**, 3501 (1963).

(7) S. J. La Placa and J. A. Ibers, *Acta Cryst.*, in press.

(8) S. J. La Placa, J. A. Ibers, and W. C. Hamilton, *J. Am. Chem. Soc.*, **86**, 2288 (1964).

(9) S. J. La Placa, W. C. Hamilton, and J. A. Ibers, *Inorg. Chem.*, **3**, 1491 (1964).

(10) W. Hieber, *Die Chemie*, **55**, 24 (1942).

(11) W. Hieber and G. Wagner, *Z. Naturforsch.*, **13b**, 339 (1958).

(12) W. E. Wilson, *ibid.*, **18b**, 349 (1958).

(13) F. A. Cotton, J. L. Down, and G. Wilkinson, *J. Chem. Soc.*, 833 (1959).

(14) E. O. Bishop, J. L. Down, P. R. Emtage, R. E. Richards, and G. Wilkinson, *ibid.*, 2484 (1959).

by Chatt, Duncanson, and Shaw.¹⁵ The crystals are air stable; of more importance, we observed none of the signs of decomposition in the X-ray beam found by Owston, *et al.*,⁴ for the related Br compound. On the basis of an optical examination and of precession and Weissenberg photographs the crystals were found to belong to the monoclinic system. The cell has the dimensions $a = 11.80 \pm 0.02$, $b = 16.93 \pm 0.03$, $c = 14.31 \pm 0.02$ Å, $\beta = 108.4 \pm 0.3^\circ$, $V = 2713$ Å³. An experimental density of 1.60 g./cm.³, obtained by flotation in zinc bromide solution, is in good agreement with a density of 1.616 g./cm.³ calculated for four molecules in the unit cell. The systematic extinctions $0k0$ for k odd and $h0l$ for l odd strongly favor the space group $C_{2h}^5-P2_1/c$. Thus all atoms are in general positions and no symmetry need be imposed on the molecule.

Intensity data were collected at room temperature by the equi-inclination Weissenberg technique, using Zr-filtered Mo K α radiation. The crystal was mounted with b as the rotation axis, and the layers $h0l$ to $h14l$ were recorded. The intensities of 1711 independent reflections accessible within the angular range $\theta_M \leq 20^\circ$ were estimated visually. After the usual Lorentz-polarization factor had been applied, the resultant F_o^2 values (where F_o is the observed structure amplitude) were corrected for absorption. For this purpose the six faces of the crystal were identified by a combination of Weissenberg photography and optical goniometry and were carefully measured. Even though the calculated volume of the crystal is only 0.00715 mm.³ and the calculated weight 11.6 μ g., the resultant transmission coefficients for a linear absorption coefficient of 56.95 cm.⁻¹ range from about 0.17 to 0.59.¹⁶ Thus the obvious fact is re-emphasized that absorption corrections are essential if one is going to achieve high accuracy in problems of this sort. The F_o values were subsequently brought to an approximate common scale through a modification of Wilson's procedure.

Solution and Refinement of the Structure

The position of the Pt atom was readily determined from a three-dimensional Patterson function. Although the positions of the Cl and P atoms undoubtedly could have been determined from the same function, it was deemed more economical in time to phase a difference Fourier with the Pt contributions. In this difference Fourier the positions of the P and Cl atoms were readily apparent. The positions of all of the carbon atoms were located in a second difference Fourier based on phases obtained from the refined positions of the heavy atoms.

The structure was refined by the least-squares procedure. The function minimized was $\sum w(F_o - F_c)^2$, where the weights w were assigned in the following way:

(15) J. Chatt, L. A. Duncanson, and B. L. Shaw, *Proc. Chem. Soc.*, 343 (1957).

(16) Programs for the IBM 7090 used in this work were local modifications of Burnham's GNABS absorption program, Zalkin's FORDAP Fourier program, and the Busing-L Levy ORFLS least-squares program, together with various local programs.

$F \leq 30 e$, $w = (F/30)^2$; $30 < F < 80 e$, $w = 1$; $F \geq 80 e$, $w = (80/F)^2$. The atomic scattering factors for the neutral atoms tabulated by Ibers¹⁷ were used. The anomalous parts of the Pt, P, and Cl scattering factors were obtained from Templeton's tabulation¹⁸ and were included in the calculated structure factors.¹⁹ Contributions of the hydrogen atoms to the structure factors were ignored. Initially the refinement was carried out with the phenyl rings constrained to their normal geometry (D_{6h} symmetry, C-C = 1.392 Å.), using the group-refinement procedure described previously.^{6,7} Each ring was assigned a single, variable isotropic thermal parameter and six variable positional parameters. These are the fractional coordinates x_o , y_o , z_o of the ring center and three angles δ , ϵ , and η which are successive counterclockwise rotations about the internal axes a_2' , a_1' , and a_3' which bring about the alignment (except for translation) of this orthogonal internal coordinate system a' with an external orthogonal coordinate system A . The origin of the internal system is taken at the ring center, with a_3' normal to the ring and a_1' intersecting a vertex. The external coordinate system has A_2 parallel to a_2 , A_1 parallel to $a_2 \times a_3$, and A_3 parallel to $A_1 \times A_2$, where a is the original monoclinic coordinate system. This initial refinement of the rings together with the other atoms, each of which was assigned a variable isotropic thermal parameter, converged rapidly to a conventional R factor ($R = \sum ||F_o| - |F_c|| / \sum |F_o|$) of 0.11 for the 60 positional and thermal parameters. A difference Fourier based on this refinement provided evidence for anisotropic thermal motion of the heavy atoms. After correction of several indexing errors, a further group refinement was carried out in which the Pt, P, and Cl atoms were assigned anisotropic thermal parameters, but the ethyl carbon atoms were assigned isotropic thermal parameters. This refinement of 80 positional and thermal parameters converged to an R factor of 0.077 and to a weighted R factor R' [$R' = (\sum w(F_o - F_c)^2 / \sum w F_o^2)^{1/2}$] of 0.098.

The obvious advantages of the group refinement are rapid convergence and minimization of computer time through the reduction of the number of variables. We have contended⁷ that an additional important advantage is that physically reasonable information can be included in the refinement, and we believe this to be preferable to imposing, for example, no constraints on the phenyl ring geometries, especially since elucidation of an unknown aspect of the structure may suffer less from correlation with errors of assumption than it would from correlation with experimental errors of measurement for the alternative full determination. In previous problems,^{7,20,21} because of their size, it has not been possible to compare the unrestricted and group refinements. In the present case, however, the un-

(17) J. A. Ibers, "International Tables for X-ray Crystallography," Kynoch Press, Birmingham, England, 1962, Vol. 3, Table 3.3.1.

(18) D. H. Templeton, *ibid.*, Table 3.3.2C.

(19) J. A. Ibers and W. C. Hamilton, *Acta Cryst.*, **17**, 781 (1964).

(20) J. A. Ibers and S. J. La Placa, *Science*, **145**, 920 (1964).

(21) S. J. La Placa and J. A. Ibers, *Inorg. Chem.*, **4**, 778 (1965).

TABLE I
 POSITIONAL, THERMAL, AND GROUP PARAMETERS FOR PtHCl[P(C₆H₅)₂C₂H₅]₂

Atom	x	y	z	β_{11}^a or B (Å. ²)	β_{22}	β_{33}	β_{12}	β_{13}	β_{23}
Pt	0.17910 (9) ^b	0.41664 (7)	0.28903 (8)	0.00423 (8)	0.00262 (5)	0.00544 (6)	-0.00045 (6)	0.00164 (6)	-0.00007 (6)
Cl	0.2588 (7)	0.4506 (6)	0.4619 (5)	0.0094 (7)	0.0067 (5)	0.0052 (5)	-0.0030 (5)	0.0019 (5)	-0.0010 (4)
P ₁	-0.0015 (6)	0.3837 (5)	0.3022 (5)	0.0059 (6)	0.0026 (4)	0.0049 (4)	-0.0003 (4)	0.0010 (4)	-0.0002 (3)
P ₂	0.3438 (5)	0.4567 (5)	0.2526 (5)	0.0040 (6)	0.0025 (4)	0.0059 (4)	0.0000 (3)	0.0014 (4)	0.0000 (3)
P ₁ EC ₁ ^c	-0.010 (2)	0.362 (2)	0.427 (2)	3.9 (6)					
P ₁ EC ₂	0.060 (3)	0.281 (3)	0.459 (3)	7.0 (9)					
P ₂ EC ₁	0.485 (2)	0.467 (2)	0.354 (2)	3.5 (5)					
P ₂ EC ₂	0.527 (3)	0.389 (2)	0.406 (2)	6.3 (8)					
Group	x_c^d	y_c	z_c	δ	ϵ	η	B (Å. ²)		
P ₁ R ₁ ^e	-0.1251 (12)	0.2284 (10)	0.1840 (9)	4.32 (1)	6.09 (1)	3.48 (1)	4.5 (3)		
P ₁ R ₂	-0.1870 (10)	0.5261 (9)	0.2273 (10)	4.39 (7)	4.87 (1)	5.28 (7)	4.8 (3)		
P ₂ R ₁	0.4306 (11)	0.3515 (10)	0.0985 (9)	2.45 (1)	0.31 (1)	5.34 (1)	4.6 (3)		
P ₂ R ₂	0.3063 (14)	0.6263 (11)	0.1468 (11)	1.90 (1)	5.93 (1)	3.62 (1)	5.8 (4)		

^a The form of the anisotropic thermal ellipsoid is $\exp[-(\beta_{11}h^2 + \beta_{22}k^2 + \beta_{33}l^2 + 2\beta_{12}hk + 2\beta_{13}hl + 2\beta_{23}kl)]$. ^b Numbers in parentheses here and in succeeding tables are estimated standard deviations in the least significant digits. ^c P₁EC₁ refers to ethyl carbon number one on P₁, etc. ^d x_c , y_c , z_c are fractional coordinates of the ring centers. The angles δ , ϵ , η (in radians) are defined in the text. ^e P₁R₁ is ring 1 on P₁, etc.

restricted refinement is feasible and for the sake of comparison it was carried out. In this calculation no geometrical restrictions were placed on the phenyl rings, and each carbon of a phenyl ring was assigned an individual, variable thermal parameter. (Again the Pt, P, and Cl atoms were assigned anisotropic thermal parameters.) This refinement of 148 positional and thermal parameters converged to an R value of 0.076 and an R' value of 0.094. It is not possible to assess directly the significance of this improvement in the usual way,²² for both the positional and thermal constraints have been removed. However, several remarks of some interest can be made from the detailed comparisons of the results of the two refinements. First, and most important, no significant shifts occurred in the parameters of the nongroup atoms; in fact, the largest shift is less than a standard deviation. Second, as expected, the distribution of the unrestricted C-C distances in the phenyl rings follows essentially a normal error curve, with C-C distances ranging from 1.30 to 1.49 Å. and averaging 1.41 ± 0.01 Å. Clearly this wide distribution, which shows no discernible trend with position in the rings, results from errors in the intensity data, rather than from forces in the crystal, and reinforces our view that if one knows geometries better than he can determine them, then this information should be put into the refinement. Third, the group refinement requires 11 min. of IBM 7094 time per cycle; the unrestricted refinement requires 24 min. per cycle.

Because of our feeling that the results of the group refinement are to be preferred on physical grounds, we give in Table I the parameter values obtained from it, together with the standard deviations of the parameters as estimated from the inverse matrix. In Table II the carbon atom positions of the phenyl rings are listed. These were derived from the data of Table I. The standard deviations in Table II are meant to be used in error analyses of functions not involved in the same ring; obviously the intra-ring distances, for example, are fixed. In Table III the values of $10F_o$ and $10|F_o|$ (in

 TABLE II
 DERIVED PARAMETERS FOR GROUP CARBON ATOMS^a

Group atom			
P ₁ R ₁ C ₁	-0.073 (2)	0.298 (1)	0.230 (1)
P ₁ R ₁ C ₂	-0.001 (1)	0.236 (1)	0.216 (1)
P ₁ R ₁ C ₃	-0.054 (2)	0.167 (1)	0.171 (1)
P ₁ R ₁ C ₄	-0.178 (2)	0.159 (1)	0.138 (1)
P ₁ R ₁ C ₅	-0.249 (1)	0.221 (1)	0.152 (1)
P ₁ R ₁ C ₆	-0.196 (2)	0.290 (1)	0.197 (1)
P ₁ R ₂ C ₁	-0.111 (2)	0.462 (1)	0.260 (1)
P ₁ R ₂ C ₂	-0.132 (2)	0.491 (1)	0.165 (1)
P ₁ R ₂ C ₃	-0.209 (2)	0.554 (1)	0.132 (1)
P ₁ R ₂ C ₄	-0.263 (2)	0.590 (1)	0.194 (1)
P ₁ R ₂ C ₅	-0.242 (2)	0.562 (1)	0.290 (1)
P ₁ R ₂ C ₆	-0.165 (2)	0.498 (1)	0.323 (1)
P ₂ R ₁ C ₁	0.393 (2)	0.398 (1)	0.164 (1)
P ₂ R ₁ C ₂	0.347 (2)	0.322 (1)	0.140 (1)
P ₂ R ₁ C ₃	0.384 (2)	0.276 (1)	0.075 (1)
P ₂ R ₁ C ₄	0.468 (2)	0.305 (1)	0.033 (1)
P ₂ R ₁ C ₅	0.514 (2)	0.381 (1)	0.057 (1)
P ₂ R ₁ C ₆	0.477 (2)	0.427 (1)	0.123 (1)
P ₂ R ₂ C ₁	0.324 (2)	0.553 (1)	0.194 (2)
P ₂ R ₂ C ₂	0.220 (2)	0.568 (1)	0.116 (2)
P ₂ R ₂ C ₃	0.202 (2)	0.641 (1)	0.070 (1)
P ₂ R ₂ C ₄	0.289 (2)	0.699 (1)	0.100 (2)
P ₂ R ₂ C ₅	0.393 (2)	0.685 (1)	0.177 (2)
P ₂ R ₂ C ₆	9.411 (2)	0.612 (1)	0.224 (1)

^a C₁ is attached to P; other C atoms are numbered in succession so that C₄ is *para* to C₁.

electrons) are listed for the 1711 reflections used in the refinement. None of the unobserved reflections accessible on the films has a calculated intensity which exceeds our estimate of the minimum observable, and hence these reflections are omitted from Table III.

The anisotropic thermal parameters can be interpreted in terms of real amplitudes of vibration, since they are essentially uncorrelated with the scale factors. (Note that in the partial anisotropic refinement the 28 carbon atoms, a reasonable fraction of the scattering power, were restricted to isotropic vibration.) In Table IV the principal values of the root-mean-square amplitudes of vibration are listed. The orientations of the thermal ellipsoids, which may be derived from

TABLE IV

ROOT-MEAN-SQUARE AMPLITUDES OF VIBRATIONS (IN Å.)			
Atom	Minimum	Intermediate	Maximum
Pt	0.157 (2)	0.199 (2)	0.227 (1)
Cl	0.20 (1)	0.23 (1)	0.34 (1)
P ₁	0.18 (1)	0.20 (1)	0.23 (1)
P ₂	0.16 (1)	0.19 (1)	0.24 (1)

the data of Table I, are not given. A rigid-body analysis has not been attempted, largely because it seems inapplicable. For example, only the Cl atom shows maximum displacement along the direction perpendicular to the Pt, P, Cl plane.

A final difference Fourier, based on the group refinement, shows peaks as high as $2.2 \text{ e}/\text{Å}^3$ in the vicinity of the Pt atom. Considerable effort was made to locate the hydride H atom, using the method of successively limiting the angular range of terms in the Fourier, a method which worked with success in our study of $\text{RhH}(\text{CO})[\text{P}(\text{C}_6\text{H}_5)_3]_3$.^{6,7} Although the noise in the vicinity of the Pt diminished considerably as the angular range was reduced, it never decreased sufficiently to enable us to discern with certainty the electron density of the hydrogen. This density is calculated to be about 0.3 to $0.5 \text{ e}/\text{Å}^3$,⁷ some 30% of the noise level on the difference Fourier with terms limited to the angular range $(\lambda^{-1} \sin \theta) \leq 0.35 \text{ Å}^{-1}$. The reasons for the noise level here are probably related to the white radiation problem always inherent with Mo $K\alpha$ radiation, to errors in the published values of μ/ρ used in the absorption correction,²³ and to an improper treatment of the platinum scattering. (The Thomas-Fermi-Dirac scattering factor for Pt used in this study differs in the range of interest by only a few tenths of an electron from the new Hartree-Fock values²⁴ and by some 1 to 2 electrons from the new Hartree-Fock-Slater values²⁵; nevertheless, the differences are smooth and should be compensated for in part by changes in the thermal parameters.) The one discernible feature on this difference map is that the phenyl carbon atoms nearest to phosphorus are vibrating less than those further away. This effect is also evident in the unrestricted refinement: the thermal parameters average on C₁ (attached to P), C₂ and C₅ (*ortho*), C₃ and C₆ (*meta*), and C₄ (*para*) 3.6, 4.6, 5.5, and 5.5 Å², respectively. Obviously one could improve the group refinement, with consequent increase in computing time, by assigning individual thermal parameters to the group atoms; this does not seem worthwhile, since the differences from the group average can always be derived from a difference Fourier.

Description of the Structure

The structure described by the cell constants, the symmetry operations of the space group, and the atomic parameters consists of the packing of discrete molecules of $\text{PtHCl}[\text{P}(\text{C}_6\text{H}_5)_2\text{C}_2\text{H}_5]_2$. The closest Pt-Pt approach is greater than 9 Å. The P atoms are *trans*

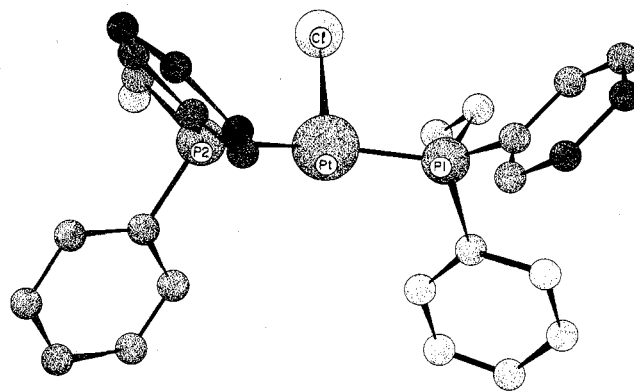


Figure 1.—A perspective drawing of one molecule of $\text{PtHCl}[\text{P}(\text{C}_6\text{H}_5)_2\text{C}_2\text{H}_5]_2$. Positions of the hydrogen atoms were not determined and are not shown.

to one another and together with the Cl atom lie at three of the four corners of a distorted square. The Pt atom is at the center of the square. This arrangement is shown in perspective in Figure 1. The important intramolecular distances are listed in Table V. The general agreement with the much less precise results of Owston, *et al.*,⁴ on the related bromine compound is exceedingly good: they found Pt-P distances of $2.26 \pm 0.07 \text{ Å}$, a Pt-Br distance of $2.56 \pm 0.04 \text{ Å}$, and Br-Pt-P angles of 94.1 and $93.7^\circ (\pm 2.0^\circ)$. The Pt-P distances of 2.267 and $2.269 \pm 0.008 \text{ Å}$ found here are toward the low end of the range found recently between P and other second or third row transition elements (Rh-P = 2.32 Å in $\text{RhH}(\text{CO})[\text{P}(\text{C}_6\text{H}_5)_3]_3$;^{6,7} Ir-P = $2.36, 2.37 \text{ Å}$ in $\text{IrO}_2\text{Cl}(\text{CO})[\text{P}(\text{C}_6\text{H}_5)_3]_2$;²⁰ Ru-P = $2.37, 2.41, 2.23 \text{ Å}$ in $\text{RuCl}_2[\text{P}(\text{C}_6\text{H}_5)_3]_3$.²¹ The Pt-Cl distance of $2.422 \pm 0.009 \text{ Å}$ is consistent with the Pt-Br distance found by Owston, *et al.*,⁴ since the covalent radius of Cl is generally considered to be some 0.12 Å less than that of Br. The Pt-Cl distance is longer than is predicted from radii-sum rules (2.30 Å), and this lengthening is consistent with the high chemical lability of the chlorine in this compound. Yet the Ru-Cl distance is 2.39 Å in $\text{RuCl}_2[\text{P}(\text{C}_6\text{H}_5)_3]_3$;²¹ and the Ir-Cl distance is approximately 2.40 Å in $\text{IrO}_2\text{Cl}(\text{CO})[\text{P}(\text{C}_6\text{H}_5)_3]_2$. Owston, *et al.*,⁴ discuss the Pt-P shortening and Pt-Br lengthening in terms of possible *trans* influence of the hydrogen atom and/or as a result of d-orbital contributions to Pt-P bonding.

TABLE V

SELECTED INTRAMOLECULAR DISTANCES AND ANGLES			
Intramolecular distance, Å.	Angles, deg.		
Pt-Cl	2.422 (9)	P ₁ -Pt-Cl	92.6 (4)
Pt-P ₁	2.267 (8)	P ₂ -Pt-Cl	94.5 (4)
Pt-P ₂	2.269 (8)	P ₁ -Pt-P ₂	188.8 (3)
Cl-P ₁	3.39 (1)	P ₁ -P ₁ EC ₁ -P ₁ EC ₂	105 (2)
Cl-P ₂	3.45 (1)	P ₂ -P ₂ EC ₁ -P ₂ EC ₂	113 (2)
P ₁ -P ₁ EC ₁	1.86 (3)	P ₁ R ₁ C ₁ -P ₁ -P ₁ EC ₁	103 (1)
P ₂ -P ₂ EC ₁	1.84 (3)	P ₁ R ₁ C ₁ -P ₁ -P ₁ R ₂ C ₁	104 (1)
P ₁ EC ₁ -P ₁ EC ₂	1.59 (5)	P ₁ R ₂ C ₁ -P ₁ -P ₁ EC ₁	102 (1)
P ₂ EC ₁ -P ₂ EC ₂	1.52 (5)	P ₂ R ₁ C ₁ -P ₂ -P ₂ EC ₁	101 (1)
P ₁ -P ₁ R ₁ C ₁	1.83 (2)	P ₂ R ₁ C ₁ -P ₂ -P ₂ R ₂ C ₁	101 (1)
P ₁ -P ₁ R ₂ C ₁	1.82 (2)	P ₂ R ₂ C ₁ -P ₂ -P ₂ EC ₁	104 (1)
P ₂ -P ₂ R ₁ C ₁	1.85 (2)		
P ₂ -P ₂ R ₂ C ₁	1.82 (2)		

(23) S. C. Abrahams, *Acta Cryst.*, **17**, 1327 (1964).

(24) D. T. Cromer, A. C. Larson, and J. T. Waber, *ibid.*, **17**, 1044 (1964).

(25) H. P. Hanson, F. Herman, J. D. Lea, and S. Skillman, *ibid.*, **17**, 1040 (1964).

Such d-orbital contributions have recently been put forth²⁶ as the explanation for the short Pt-As bond lengths of 2.38 Å. in Pt(diarsine)₂I₂²⁷ and in Pt(diarsine)₂Cl₂.²⁶ Since the covalent radius of As is some 0.10 Å. longer than that of P, the contraction of the Pt-As bonds is similar to that of the Pt-P bonds found here. Nevertheless, there is sufficiently little information on M-P and M-Cl distances (where M is a second or third row transition element) so that very little can be said with certainty about the bonding in this platinum compound. It is not possible to assess the relative importance of the *trans* effect or of d-orbital contributions, or in fact of any other effects that one might fancy. Yet, on the basis of the spreading of the P-Pt-Cl angles from the presumably ideal value of 90° to 92.6 and 94.5° there is the indication of some steric interaction between Cl and P. (A similar deviation of the angles from 90° was observed in the bromo compound.⁴) Hence any factor that tends to shorten the Pt-P bond distance will, through steric effects, lengthen the Pt-Cl distance and *vice versa*.

(26) N. C. Stephenson, *Acta Cryst.*, **17**, 1517 (1964).(27) N. C. Stephenson, *J. Inorg. Nucl. Chem.*, **24**, 791 (1963).

Of some interest in this regard is that the Pt, P, Cl portion of the molecule is significantly nonplanar. The best least-squares plane²⁸ through these four atoms has the equation $3.528x - 15.805y + 1.336z = 5.572$ (monoclinic coordinates). The deviations from this plane are Pt, 0.005 ± 0.001 ; Cl, -0.021 ± 0.010 ; P₁, -0.095 ± 0.008 ; and P₂, -0.096 ± 0.008 Å. A careful examination of the intermolecular distances suggests no explanation for a lack of planarity in terms of packing distortions. (In fact all intermolecular contacts appear to be normal, and hence are not tabulated here.) The only explanation for this lack of planarity that we can suggest involves the intramolecular steric repulsions of P and Cl.

The geometry of the diphenylethylphosphine ligand closely resembles that found previously for the triphenylphosphine ligand.^{7,20,21}

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(28) W. C. Hamilton, *Acta Cryst.*, **14**, 185 (1961).

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A Five-Coordinated d⁶ Complex: Structure of Dichlorotris(triphenylphosphine)ruthenium(II)¹

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A complete X-ray structure determination of dichlorotris(triphenylphosphine)ruthenium(II), RuCl₂[P(C₆H₅)₃]₃, has been carried out in order to ascertain if in this compound the Ru(II) (d⁶) is five-coordinated. The compound crystallizes in space group C_{2h}⁵-P2₁/c of the monoclinic system in a cell of dimensions $a = 18.01$, $b = 20.22$, $c = 12.36$ Å, $\beta = 90.5^\circ$. The observed density is 1.43 g./cm.³, and the density calculated for four molecules in the cell is 1.415 g./cm.³. There are no crystallographic symmetry conditions imposed on the molecule. The structure consists of the packing of individual monomeric molecules. In these molecules the Ru lies toward the center of a distorted square pyramid which consists of *trans* Cl atoms and *trans* P atoms in the base and an apical P atom. The apical Ru-P distance of 2.23 Å. is about 0.16 Å. shorter than the basal Ru-P distances. This is the only known example from X-ray structural studies of a five-coordinated d⁶ complex, and it probably occurs not because of an inherent tendency toward five-coordination, but rather because the unused octahedral site about the square-pyramidal configuration is effectively blocked by a phenyl ring.

Introduction

Of the two dozen or so five-coordinated transition metal complexes that have been established by X-ray diffraction studies, none is a d⁶ complex.² This is not surprising, for it has frequently been noted³ that the spin-paired d⁶ configuration is an especially favorable one for the formation of octahedral complexes. Of the relatively few d⁶ complexes thought to be five-coordinated that are reported in the literature, the case for five-coordinated Os(II) in dibromotris(tri-

phenylphosphine)osmium(II) seems one of the most convincing. This compound was prepared by Vaska⁴ in 90% yield from the reaction of (NH₄)₂OsBr₆ with triphenylphosphine in 2-methoxyethanol at 25°. In a similar manner Vaska⁵ has prepared dichlorotris(triphenylphosphine)ruthenium(II). From magnetic, molecular weight, conductivity, and spectroscopic measurements Vaska concludes that these compounds are diamagnetic, monomeric, and are not hydrides. A preliminary X-ray examination by Pollack⁶ indicated that the

(1) Research performed under the auspices of the U. S. Atomic Energy Commission.

(2) J. A. Ibers, *Ann. Rev. Phys. Chem.*, in press.(3) For example, see R. S. Nyholm, *Proc. Chem. Soc.*, 273 (1961).(4) L. Vaska, *Chem. Ind. (London)*, 1402 (1963).

(5) L. Vaska, private communication.

(6) S. S. Pollack, as quoted by Vaska.⁵